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Original research article

# Performance analysis of weather research and forecasting model for simulating near-surface optical turbulence over land

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## ABSTRACT

Accurate knowledge of the optical turbulence ( $C_n^2$ ) in the atmospheric surface layer is of great significance for the applications of various ground-based optical systems (e.g., astronomical observation, laser communication and target detection), however, many methods using instruments to measure  $C_n^2$  are limited spatially and temporally and they are more difficult and expensive in severe environments, which makes estimations a less expensive and convenient alternative. This study introduces a method where the products of Weather Research and Forecasting (WRF) model coupled with Monin–Obukhov similarity (MOS) theory is used to estimate the value of  $C_n^2$  in the surface layer, and the performance of this method is evaluated against the corresponding  $C_n^2$  dataset from the *in-situ* micro-thermometer during a campaign over south-west China. The variance between the estimated  $C_n^2$  and measured  $C_n^2$  is found to be comparable to the observed variance of the temporal due to temporal and spatial variations of the atmospheric turbulence. Evaluation against direct observations shows that this approach can well estimate  $C_n^2$  time series with some error in some time periods, and both the estimated  $C_n^2$  and the measured  $C_n^2$  are more consistent in trend and magnitude in general. Given these validated results, this method has the potential to be a useful tool in a wide range of ground-based optical applications.

## 1. Introduction

Atmospheric turbulence intensity is usually characterized by the refractive index structure constant ( $C_n^2$ ,  $\text{m}^{-2/3}$ ), and various parameters with  $C_n^2$  as the integral parameter such as the seeing ( $\epsilon$ ), the isoplanatic angle ( $\theta_0$ ) and the wavefront coherence time ( $\tau_0$ ) are important for studying the effect of laser transmission in the atmosphere [1–3].

Various instruments such as Scintillometer [4], Interferometers [5], Differential Image Motion Monitor (DIMM) [6], Micro-thermometry (Balloon) [7], Scintillation Detection and Ranging (SCIDAR) [8], Multi Aperture Scintillation Sensor (MASS) [9] have been used to measure these atmospheric turbulence parameters. These instruments are not easy to operate and are quite expensive, thus, it makes the instruments difficult to perform with a large-scale, long-term, systematic observation, and also in strong winds and freezing conditions they seem too fragile to obtain reliable results. Therefore, there is a need for a robust method for estimating  $C_n^2$  based on readily routine meteorological parameters (e.g., temperature, relative humidity, pressure, wind speed, etc.).

Monin–Obukhov similarity (MOS) theory provides an important theoretical basis to obtain  $C_n^2$  from routine meteorological parameters [10–15]. Over the years, Masciadri and Lascaux have used Mesoscale Non-hydrostatic (Meso-Nh) model to obtain  $C_n^2$  profile above the sites of interest for astronomy (e.g., Roque de Los Muchachos Observatory; European Southern Observatory)

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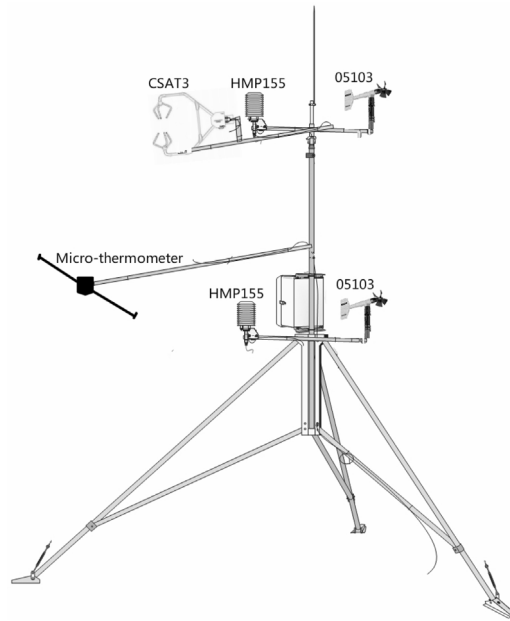


Fig. 1. The *in-situ* measurement system.

[16–21]. However, using Weather Research and Forecasting (WRF) model and Monin–Obukhov similarity (MOS) theory to estimate  $C_n^2$  in the surface layer is seldom reported. On account of the various ground based optical application requirements, the purpose of this paper is to present a significant result of estimated  $C_n^2$  in the surface layer over land, and this method can be used for gap filling in case of failures of the instruments.

The remainder of this paper is organized as follows. The detailed description of *in-situ* measurement and the configuration of modeling are given in Section 2. The methodologies of estimating  $C_n^2$  is introduced in Section 3. The estimation  $C_n^2$  results and the related statistical analysis are shown in Section 4. Finally, conclusions and discussion are given in Section 5.

## 2. Data and methodologies

### 2.1. The description of observational site and environment

The measurement place is located in Chengdu area (south-west China) in which belongs to a subtropical wet monsoon climate. Because of its topography and other special geographical and environmental factors, it also has a prominent vertical climate. The *in-situ* measurement system includes a data logger (CR5000, details specifications are given in the manufacturer <https://campbell.secure.force.com/campbellsite>), ultrasonic anemometer (CSAT3), micro-thermometer, temperature and relative humidity probe (HMP155), wind speed & direction sensors (05103), CS106 barometer, 485 communication module, power module and a 3-m tower, which are shown in Fig. 1.

HMP155 sensor and 05103 sensor are installed on the first floor, approximately 1 m above the ground. CSAT3 ultrasonic anemometers and micro-thermometer are installed on the second floor approximately 3 m above the ground. The underlying surface of observational site is low shrubbery and grassland, but there is a house about 60 m high and 500 m away from the measurement site.

Table 1

Main technical specification of the *in-situ* measurement system. Note that temperature, relative humidity, wind speed, wind direction and pressure are abbreviated as T, RH, WS, WD and P in the table.

Instrument	Parameter	Range	Accuracy
Micro-thermometer	$C_n^2$	$10^{-18} \sim 10^{-12} \text{ m}^{-2/3}$	
CSAT3	WS	$u_x: \pm 30 \text{ m s}^{-1}$ $u_y: \pm 30 \text{ m s}^{-1}$ $u_z: \pm 8 \text{ m s}^{-1}$	$u_x \leq \pm 4 \text{ cm s}^{-1}$ $u_y \leq \pm 4 \text{ cm s}^{-1}$ $u_z \leq \pm 2 \text{ cm s}^{-1}$
HMP155 sensor	T	$-40 \sim +60 \text{ }^\circ\text{C}$	$\leq \pm 0.5^\circ\text{C}$
	RH	$0 \sim 100\%$	$\pm 1\%$
05103 sensor	WS	$0 \sim 100 \text{ m s}^{-1}$	$0.5 \text{ m s}^{-1}$
	WD	$0 \sim 360^\circ$	$3^\circ$
CS106 barometer	P	$500 \sim 1100 \text{ hPa}$	$\pm 0.3 \text{ hPa @ } 20^\circ\text{C}$

The main technical indicators of the *in-situ* measurement system are shown in Table 1.

2.2.  $C_n^2$  data from the micro-thermometer

For visible and near-infrared wavelengths light,  $C_n^2$  is mainly caused by temperature fluctuation. In this study,  $C_n^2$  values deduced from the micro-thermometer are used for comparison, and they are measured by a pair of horizontally micro-temperature probes. Temperature structure function  $D_T(r)$  is one of statistics describing temperature disturbance at two points separated by distance  $r$  [22], and is given by

$$D_T(r) = \langle [T(\vec{x}) - T(\vec{x} + \vec{r})]^2 \rangle, \quad l_0 \ll r \ll L_0, \tag{1}$$

where  $T$  is air temperature (K),  $\vec{x}$  and  $\vec{r}$  denote the position vector,  $r$  is the magnitude of  $\vec{r}$ ,  $\langle \dots \rangle$  represents ensemble average,  $l_0$  and  $L_0$  are the inner and outer scales of atmospheric turbulence. When  $r$  is between  $l_0$  and  $L_0$ , this function appears as a power law in  $r^{2/3}$ . These scales can be seen as the minimum and maximum sizes of turbulent eddies. In the inertia range  $[l_0, L_0]$ , defines the structure constant of temperature ( $C_T^2$ ) as below

$$D_T(r) = C_T^2 r^{2/3}, \quad l_0 \ll r \ll L_0. \tag{2}$$

The relationship between  $C_n^2$  and  $C_T^2$  is expressed by [23] and using the Gladstone's law that relates the refractive index to temperature structure constant,  $C_n^2$  is expressed as

$$C_n^2 = \left( 79 \times 10^{-6} \frac{P}{T^2} \right)^2 C_T^2, \tag{3}$$

where  $P$  is air pressure (hPa).

The two probes are legs of a Wheatstone bridge, and the resistance of probe is nearly proportional to temperature, thus, the temperature change can be transformed into the resistance change to obtain the voltage change. The micro-thermometer system provides  $C_T^2$  data by measuring mean square temperature fluctuations from Eq. (1) and thus  $C_n^2$  value can be acquired from Eqs. (2)–(3). In our case, a 10  $\mu\text{m}$  diameter platinum wire has a standard deviation of minimum temperature fluctuation is less than 0.002 K [24].

In order to verify the performance of this method, 6 days of available  $C_n^2$  observations data measured by micro-thermometer are used to as a comparison. The average routine meteorological parameters and weather conditions during the period of experiment are listed in Table 2.

2.3. Model configuration

WRF model is a mesoscale atmospheric model, which is developed by the National Centers for Environmental Prediction (NCEP) and the National Center for Atmospheric Research (NCAR) of the United States. Detailed control equations, transformations and grid adaptation, see the User's Guide [25]. In this study, the Advance Research WRF (ARW) model is used to conduct simulations.

There are 40 vertical levels and three nested horizontal domains with horizontal resolutions of 9 km, 3 km, and 1 km, respectively, in each simulation. The center grid point (30.57°N, 104.02°E) of the finest nested grid is close to *in-situ* measurement system. The final operational global analysis (FNL) data with 1° × 1° horizontal resolution and 6-h temporal resolution is used to initialize the input to the WRF model, which is downloaded from the web site of NCEP. The outputs are saved as 10 min intervals, and the results from first 24 h are excluded to avoid errors. WRF model relies on a selected physical schemes to output multiple meteorological parameters, see User's Guide to gain a better understanding of mesoscale modeling, and model configurations and main physical scheme options are summarized in Table 3.

3. The methodologies of estimating  $C_n^2$

According to Monin-Obukhov similarity theory,  $C_n^2$  can be defined in terms of the temperature structure parameter ( $C_T^2$ ), humidity structure parameter ( $C_q^2$ ) and temperature-humidity structure parameter ( $C_{Tq}$ ) as follows [34]:

Table 2

Average routine meteorological parameters and weather conditions during the period of experiment. Note that temperature, relative humidity, wind speed and pressure are abbreviated as T, RH, WS and P in the table.

Date	T (°C)	RH (%)	P (hPa)	WS (ms <sup>-1</sup> )	Weather condition
2014-5-1	19.5	54	956	2.2	Cloudy
2014-5-2	18.0	55	961	1.7	Overcast
2014-5-6	20.0	45	957	2.0	Sunny
2014-5-7	21.5	50	954	2.3	Sunny
2014-5-12	22.5	49	951	2.0	Cloudy
2014-5-13	22.0	52	948	1.9	Cloudy

**Table 3**  
WRF model configurations.

Attribute	Model Configuration
Simulation period	1–13 May 2014
Domain	Nested regional domains
Center longitude and latitude	104.02°E, 30.57°N
Horizontal resolution, grid point numbers	9 km, 90 × 90 (Coarse Grid) 3 km, 75 × 75 (Fine Grid) 1 km, 45 × 45 (Finest Grid)
Vertical resolution	40 layers from 1000-10 hPa
Initialized meteorological data	NCEP-FNL
Microphysics	WSM-5 [26]
Longwave Radiation	RRTM [27]
Shortwave Radiation	Goddard [28]
Surface Layer	Eta similarity [29]
Land Surface	NOAH [30,31]
Planetary boundary layer	MYJ [32,33]

$$C_n^2 = A^2 C_T^2 + 2ABC_{Tq} + B^2 C_q^2 . \tag{4}$$

At a interest wavelength of 0.55 μm,  $A = 79.0 \times 10^{-6} \frac{p}{T^2}$  and  $B = -56.4 \times 10^{-6}$ . The structure constant  $C_T^2$ ,  $C_q^2$  and  $C_{Tq}$  can be expressed as follows:

$$C_T^2 = T_*^2 z^{-2/3} f_T(z/L) , \tag{5a}$$

$$C_q^2 = q_*^2 z^{-2/3} f_q(z/L) , \tag{5b}$$

$$C_{Tq} = \gamma_{Tq} T_* q_* z^{-2/3} f_{Tq}(z/L) , \tag{5c}$$

where  $z$  is the height above the surface,  $L$  is the Obukhov length.  $T_*$  and  $q_*$  are the scaling parameters for temperature, specific humidity, respectively.  $\gamma_{Tq}$  is the correlation coefficient. The similarity functions  $f_T(z/L)$ ,  $f_q(z/L)$  and  $f_{Tq}(z/L)$  are experimentally determined, and the well-known similarity function is given by [10]:

$$f(z/L) = \begin{cases} 4.9(1 - 7z/L)^{-2/3}, & z/L < 0, \\ 4.9[(1 + 2.4(z/L)^{2/3})], & z/L > 0. \end{cases} \tag{6}$$

Subsequently, substituting Eq. (5a) into Eq. (4) gives a  $C_n^2$  expression in terms of  $T_*$ ,  $q_*$  and  $L$  as below:

$$C_n^2 = z^{-2/3} f(z/L) (A^2 T_*^2 + 2AB\gamma_{Tq} T_* q_* + B^2 q_*^2). \tag{7}$$

From Monin-Obukhov similarity theory, conditions are assumed to be horizontally homogeneous and stationary, and it is possible to define  $T_*$ ,  $q_*$  and  $u_*$  in terms of mean horizontal-vertical velocity covariance ( $\omega u$ ), mean vertical velocity-temperature covariance ( $\omega T$ ) and mean vertical velocity-water vapour density covariance ( $\omega q$ ) [35] as follows:

$$u_*^2 = -\omega u , \tag{8a}$$

$$u_* T_* = -\omega T , \tag{8b}$$

$$u_* q_* = -\omega q \tag{8c}$$

where  $u$  is the horizontal component of turbulent velocity vector, and  $\omega$  is the vertical component of turbulent velocity vector.

The bulk aerodynamic formulas are simple empirical parameterizations of the surface fluxes in the surface layer in terms of average surface layer variables, which are given by:

$$u_*^2 = C_D \bar{U}^2 , \tag{9a}$$

$$\omega T = C_H \bar{U} [T_s - T(z)] = C_H \bar{U} \Delta T , \tag{9b}$$

$$\omega q = C_E \bar{U} [q_s - q(z)] = C_E \bar{U} \Delta q , \tag{9c}$$

where  $\bar{U}$  is the mean wind speed at height  $z$ , and  $\Delta T = T_s - T(z)$ ,  $\Delta q = q_s - q(z)$ .  $T(z)$  and  $q(z)$  are mean air temperature, mean specific humidity at height  $z$ , respectively.  $T_s$  and  $q_s$  are the surface temperature, specific humidity, respectively.  $C_D$ ,  $C_H$  and  $C_E$  are the drag coefficient, sensible heat flux coefficient and moisture flux coefficient, respectively.

Consequently, substituting Eq. (9a) into Eq. (8a) gives the expression  $T_*$ ,  $q_*$  and  $u_*$  as follow:

$$u_*^2 = C_D \bar{U}^2 , \tag{10a}$$

$$T_* = -\frac{C_H \bar{U} \Delta T}{u_*}, \tag{10b}$$

$$q_* = -\frac{C_E \bar{U} \Delta q}{u_*}. \tag{10c}$$

The drag coefficient ( $C_D$ ), sensible heat flux coefficient ( $C_H$ ) and moisture flux coefficient ( $C_E$ ) are all from experiment [36]:

$$C_D = 1.3 \times 10^{-3}, \tag{11a}$$

$$C_H = \begin{cases} 0.91 \times 10^{-3}, & -20 < \bar{U} \Delta T < 25 \text{ (mKs}^{-1}\text{)}, \\ 1.46 \times 10^{-3}, & \bar{U} \Delta T > 25 \text{ (mKs}^{-1}\text{)}. \end{cases} \tag{11b}$$

$$C_E = 1.32 \times 10^{-3}. \tag{11c}$$

For the parameterization of Obukhov length  $L$ , the expression of  $L$  is given by two  $\bar{U} \Delta T$  ranges as follow:

$$L^{-1} = \begin{cases} -0.231 \left( \frac{\Delta T}{\bar{U}^2} + 0.212 \frac{\Delta q}{\bar{U}^2} \right), & -20 < \bar{U} \Delta T < 25 \text{ (mKsec}^{-1}\text{)}, \\ -0.371 \left( \frac{\Delta T}{\bar{U}^2} + 0.132 \frac{\Delta q}{\bar{U}^2} \right), & \bar{U} \Delta T > 25 \text{ (mKsec}^{-1}\text{)}. \end{cases} \tag{12}$$

Therefore, the values of  $U(z)$ ,  $T(z)$  and  $q(z)$ , as well as  $T_s$  and  $q_s$  are exported by WRF model first, then air-surface difference for temperature ( $\Delta T$ ), specific humidity ( $\Delta q$ ) and mean wind speed ( $\bar{U}$ ) are calculated.  $u_*$ ,  $T_*$ ,  $q_*$  and  $L$  can be obtained by solving Eqs. (10a)–(12). Finally, it is simple to calculate  $C_n^2$  from Eqs. (6)–(7).

### 4. Results

Fig. 2 shows the  $C_n^2$  time-series between estimated by WRF model and *in-situ* measured by micro-thermometer in the surface layer. For ease of comparison, the measured values are also superimposed on the figure. It can be seen that both the estimated  $C_n^2$  and that measured by micro-thermometer are more consistent in trend and magnitude in general. The estimated  $C_n^2$  values tend to be underestimated in the nighttime, however, in the strong turbulence are more accurate.

As we all know that Wyngaard proposed a similarity function model to estimate  $C_n^2$  in the atmospheric surface layer [10], and the

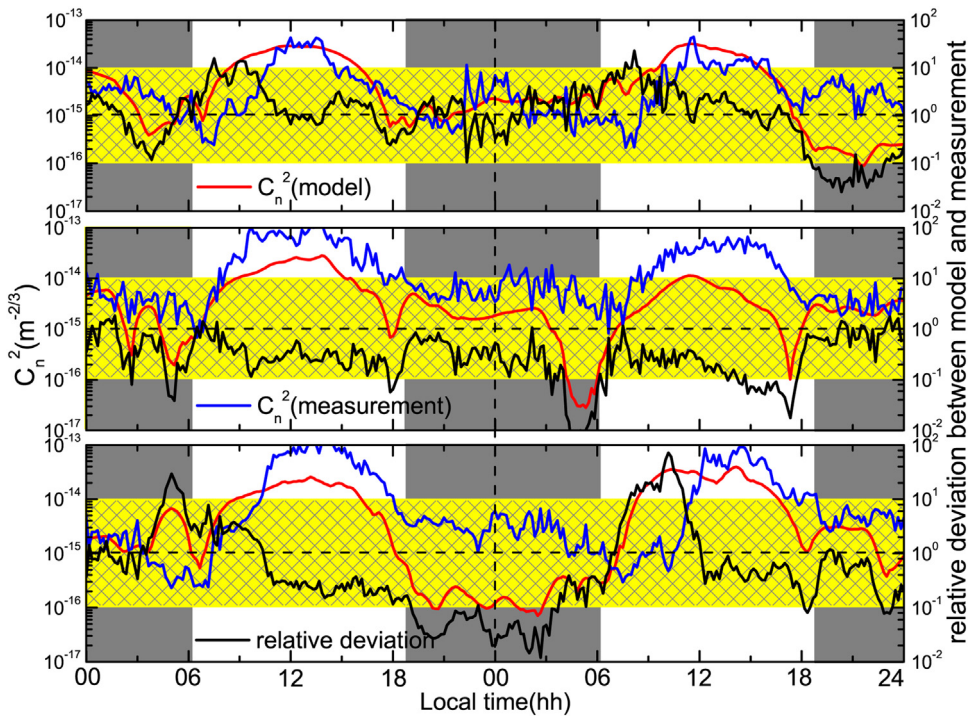


Fig. 2. The time-series of  $C_n^2$  in the surface layer. The red line and blue line represent the estimation and measurement, and the black line represents the relative deviation (right ordinate), respectively. The gray shade indicates the nighttime, and the yellow shaded area denotes the relative deviation in the range [0.1, 10]. Note that local time = UTC + 8h.

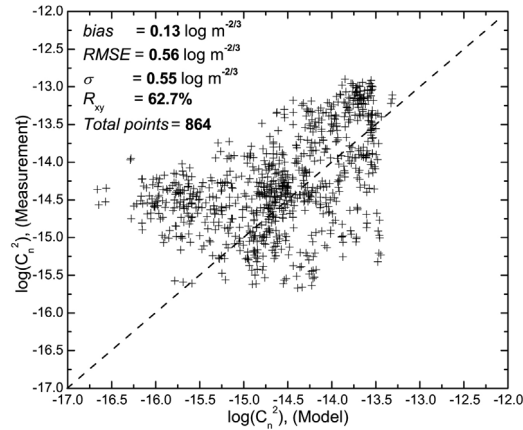


Fig. 3. Scatter diagrams of  $C_n^2$  between estimation and measurement, and each cross point represents an average of 10 min resolution.

performance of this similarity function model is satisfactory well under unstable conditions, but it is not very good under stable conditions. Literature pointed out that this similarity function model only covers a relatively short atmospheric stability range, and thus significantly limits its applicability at night under stable conditions [37]. Therefore, there is a need for constructing a new similarity function model which covers a wide range of atmospheric stabilities conditions in the future work. From Fig. 2, the fluctuations of relative deviation between the estimated  $C_n^2$  and the measured  $C_n^2$  are generally within 10 times, and the relative deviation between the two will be greater at night and the atmosphere is under stable conditions.

Three statistical operators: *bias*, root mean square error (*RMSE*), bias-corrected *RMSE* ( $\sigma$ ), and the correlation coefficient ( $R_{xy}$ ) are used to evaluate the accuracy of estimated  $C_n^2$ , which follows the same framework as [21]. For convenience,  $C_n^2$  is plotted on the logarithmic scale. The correlation of the  $\log(C_n^2)$  between estimation and measurement is depicted in Fig. 3, one can see that the *bias*, *RMSE* and  $\sigma$  are relatively small, and  $R_{xy}$  is relatively large, showing the estimation is coherent with the measurement well.

In addition, a  $3 \times 3$  contingency tables are used to study the relationship between model values and measured values, referring to the paper for details [21]. A contingency table allows for analysis of the relationship between two or more categorical variables, which is a table with  $n \times n$  entries that displays the distribution of model and measurement in terms of frequencies or relative frequencies. An example of a  $3 \times 3$  contingency table is shown in Table 4, and there are two cases in which a “hit” representing the estimated value is correct, yet a “miss” representing the estimated value is incorrect.

Using  $a, b, c, d, e, f, g, h, i$  and  $N$  ( $N = a + b + c + d + e + f + g + h + i$ ), some probabilities parameters can be calculated which are used to investigate how well (or bad) this method performs for a particular parameter. Some evaluation parameters we will use in the following paper from the generic  $3 \times 3$  contingency table of Table 4: percent of correct detections (*PC*, in %), probability of detection (*POD*, in %) and extremely bad detection (*EBD*, in %) are listed. In a random estimation case, all *POD* are equal to 33%,  $PC = 33\%$ ,  $EBD = 22.2\%$ , and the estimation will be useful if these values (*PC*, *POD*, *EBD*) perform better than those random cases. These parameters are a good reference to evaluate the performances of this approach, and we will write  $POD_i$  instead of *POD* (*event i*) with the *i* event considered.

$$PC = \frac{a + e + i}{N} \times 100, \quad 0\% \leq PC \leq 100\%, \tag{13a}$$

$$POD(\text{event 1}) = \frac{a}{a + d + g} \times 100, \quad 0\% \leq POD \leq 100\%, \tag{13b}$$

$$POD(\text{event 2}) = \frac{e}{b + e + h} \times 100, \quad 0\% \leq POD \leq 100\%, \tag{13c}$$

$$POD(\text{event 3}) = \frac{i}{c + f + i} \times 100, \quad 0\% \leq POD \leq 100\%, \tag{13d}$$

Table 4  
A generic  $3 \times 3$  contingency table.

	Intervals	Measurement			Total
		1	2	3	
Model	1	<i>a</i> (hit 1)	<i>b</i>	<i>c</i> (miss)	$a + b + c$
	2	<i>d</i>	<i>e</i> (hit 2)	<i>f</i>	$d + e + f$
	3	<i>g</i> (miss)	<i>h</i>	<i>i</i> (hit 3)	$g + h + i$
	Total	$a + d + g$	$b + e + h$	$c + f + i$	$N$

**Table 5**

3 × 3 contingency table of  $\log(C_n^2)$  ( $\text{m}^{-2/3}$ ) for WRF model and micro-thermometer in surface layer, this two thresholds -14.7 and -14.2 are defined with the climatological tertiles. Note that  $PC = 50.1\%$ ,  $EBD = 8.1\%$ ,  $POD_1 = 48.1\%$ ,  $POD_2 = 37.5\%$ ,  $POD_3 = 63.7\%$ .

	Intervals	Measurement			Total
		[, -14.7]	[-14.7, -14.2]	[-14.2,]	
Model	[, -14.7]	88	169	40	297
	[-14.7, -14.2]	65	127	84	276
	[-14.2,]	30	43	218	291
	Total	183	339	342	864

$$EBD = \frac{c + g}{N} \times 100, \quad 0\% \leq EBD \leq 100\% . \quad (13e)$$

In Table 5, the  $PC$  (50.1%) is significantly better than 33%. Moreover, the  $EBD$  (8.1%) is small which is the sign that this approach seldom produces extremely bad simulation.  $POD_1$  is 48.1%,  $POD_2$  is 37.5% and  $POD_3$  is 63.7%. In all cases,  $PC$  and  $POD_{i=1,2,3}$  are well larger than 33%.

## 5. Conclusions and discussion

The main result of this paper is that WRF model coupling MOS theory is used to rebuild the time-series of  $C_n^2$  in the surface layer over southwest China, and the results of this study provide an alternative way to obtain  $C_n^2$  when using optoelectronic systems in some scenes that are difficult to measure. Analysis shows that the  $C_n^2$  values between estimation and measurement are more consistent in trend and magnitude in general, however, the estimated  $C_n^2$  values tend to be underestimated in the nighttime. It should be pointed out that the similarity function model (W73, [10]) we use in this study only covers a relatively short atmospheric limited range, therefore, there is a need for constructing a new similarity function model in the future work. Furthermore, MOS theory makes an assumption of horizontal homogeneity, and the sensible heat flux and latent heat flux are really small in the nighttime, thus the sensible heat flux and latent heat flux yielded from WRF model output which may involve more complex error source. To be certain, more validation studies would be carried out in next work.

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